

Multiresolution Coupled-Cluster

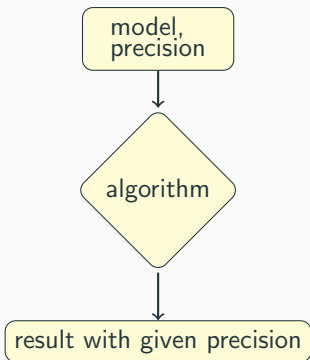
Jakob S. Kottmann

25.10.2018

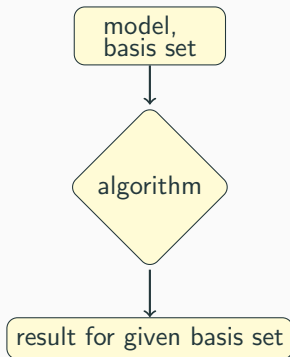
Theoretische Chemie
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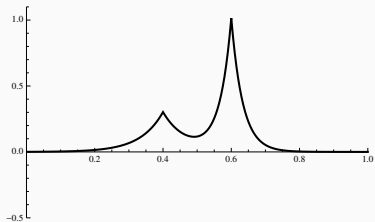
Multiresolution



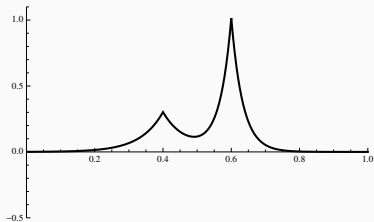
Fixed-Basis



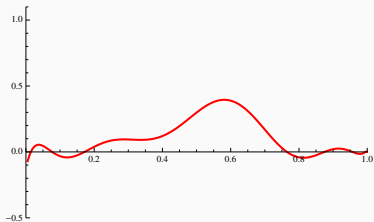
MRA - Multiresolution Analysis



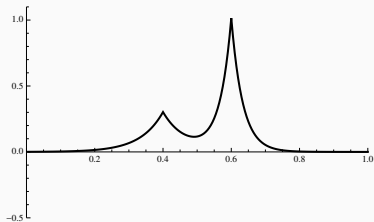
polynomials as universal basis



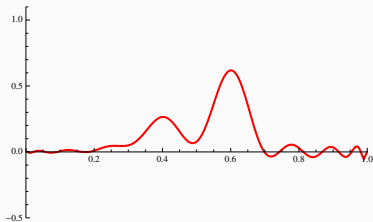
10 basis functions



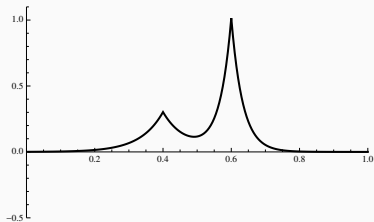
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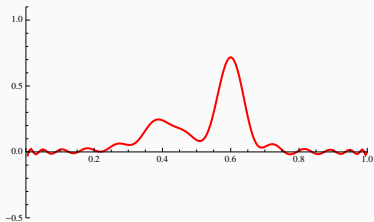
20 basis functions



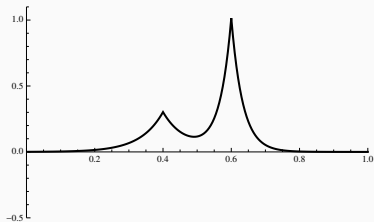
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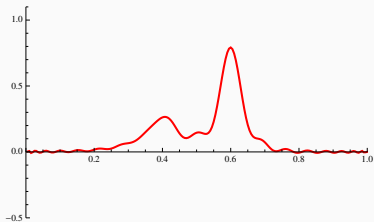
30 basis functions



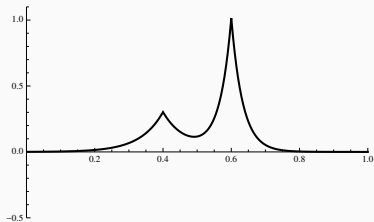
polynomials as universal basis



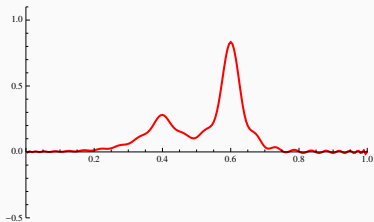
40 basis functions



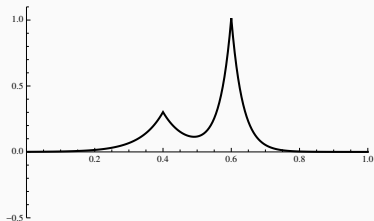
polynomials as universal basis



50 basis functions

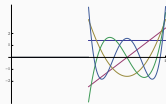
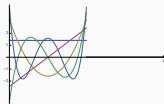
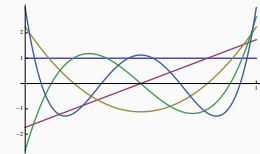


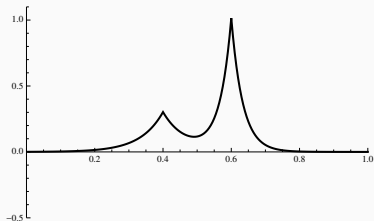
polynomials as universal basis



polynomials as universal basis

or piecewise polynomials

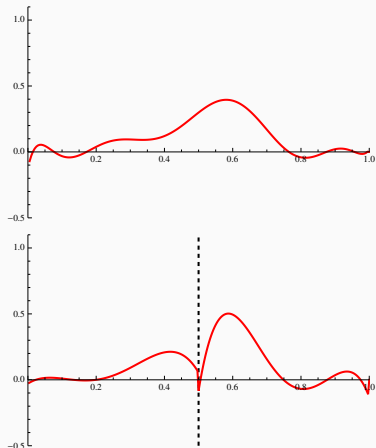


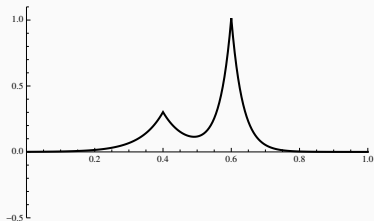


polynomials as universal basis

or piecewise polynomials

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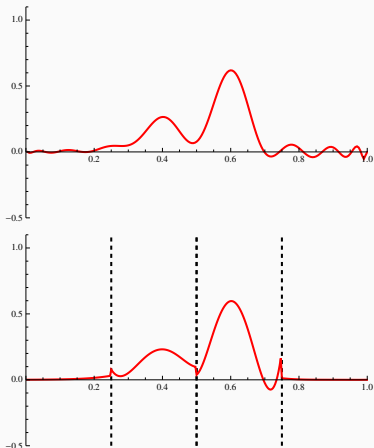


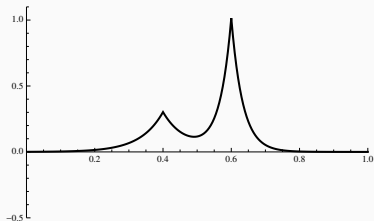


polynomials as universal basis

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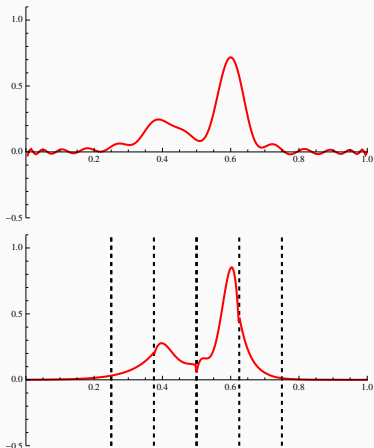


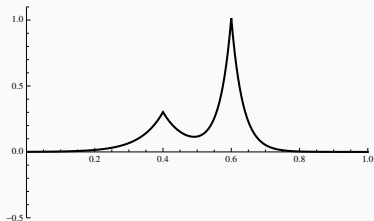


polynomials as universal basis

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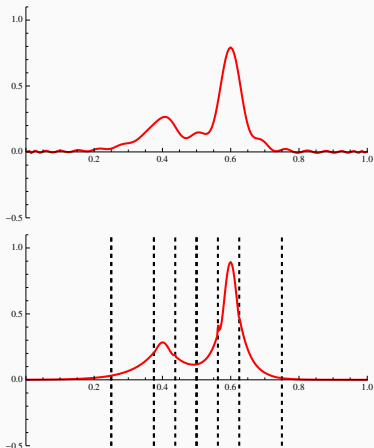


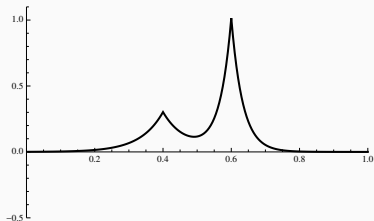


polynomials as universal basis

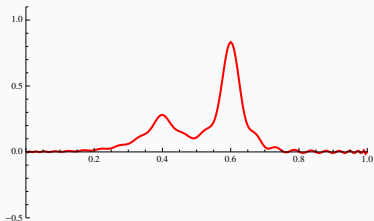
or piecewise polynomials

40 basis functions



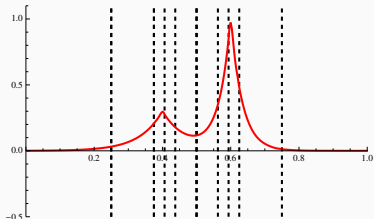


50 basis functions



polynomials as universal basis

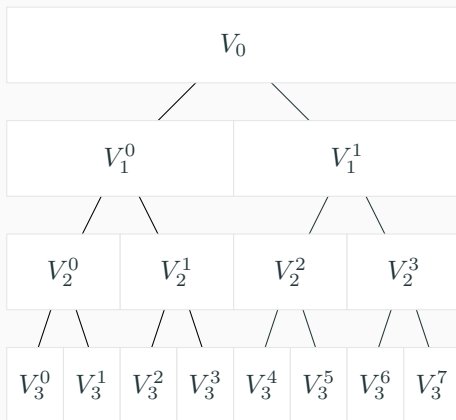
or piecewise polynomials



$$V_0 \subset V_1 \subset V_2 \subset \dots \subset \mathbf{L}^2$$

$$V_n = \bigoplus_{l=0}^{2^n - 1} V_n^l$$

$$V_{n+1} = V_n \oplus W_n$$



$$V_0 \subset V_1 \subset V_2 \subset \dots \subset L^2$$

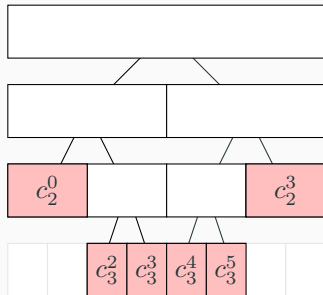
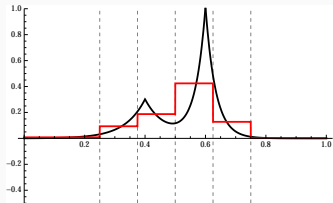
Haar Wavelets

$$V_n = \bigoplus_{l=0}^{2^n-1} V_n^l$$

$$V_1 = \left(\begin{array}{|c|} \hline \text{[Red step function: 1 on } [0, 1/2], 0 \text{ elsewhere]} \\ \hline \end{array} , \begin{array}{|c|} \hline \text{[Red step function: 1 on } [1/2, 1], 0 \text{ elsewhere]} \\ \hline \end{array} \right)$$

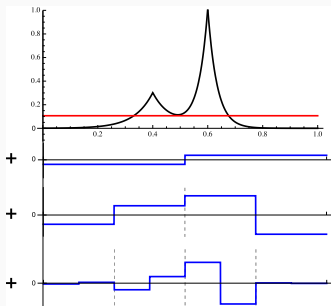
$$V_{n+1} = V_n \oplus W_n$$

$$V_1 = \left(\begin{array}{|c|} \hline \text{[Red step function: 1 on } [0, 1], 0 \text{ elsewhere]} \\ \hline \end{array} , \begin{array}{|c|} \hline \text{[Blue step function: 1 on } [0, 1/2], -1 \text{ on } [1/2, 1], 0 \text{ elsewhere]} \\ \hline \end{array} \right)$$



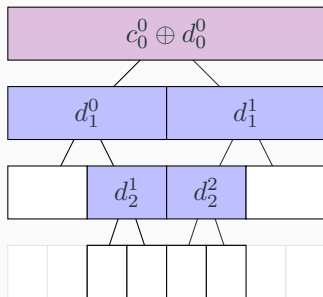
$$|f\rangle = \sum_{nl} c_n^l |\varphi_n^l\rangle$$

reconstructed representation

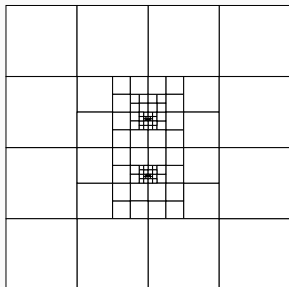
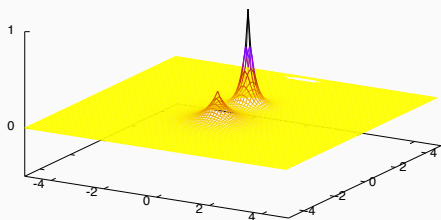


$$|f\rangle = c_0|\varphi_0\rangle + \sum_{nl} d_n^l|\psi_n^l\rangle$$

$$\varphi_n^l \in V_n^l, \quad \psi_n^l \in W_n^l$$



compressed representation



N -dimensional MRA with k scaling functions :

2^N children for each box \rightarrow Regularization

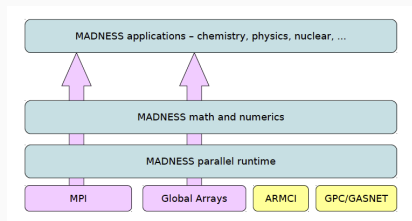
k^N coefficients for each box \rightarrow Low-Rank decomposition

$$\left(-\frac{\Delta}{2} + \hat{V}\right) |\Psi\rangle = E|\Psi\rangle$$

$$\Psi(x) = -2 \int dx G(x-x') \left(\hat{V}\Psi\right)(x')$$

$$(\Delta - 2E) G(x) = \delta(x)$$

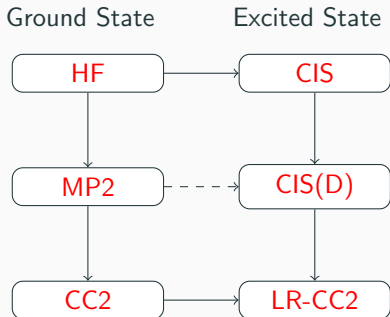
MADNESS (www.github.com/m-a-d-n-e-s-s)¹



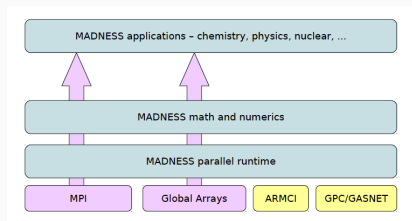
Lecture notes online²

¹R. J. Harrison *et al.*, *SIAM J. Sci. Comput.*,38:5, 2016.

²<https://sites.google.com/view/numericalquantumchemistry2018/lecture-notes>



MADNESS (www.github.com/m-a-d-n-e-s-s)¹



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¹R. J. Harrison *et al.*, *SIAM J. Sci. Comput.*,38:5, 2016.

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Real-Space Coupled-Cluster



$$\mathcal{H}|0\rangle = E|0\rangle$$

Reference Wavefunction $|0\rangle$

$$\mathcal{H} = e^{-\hat{\mathcal{T}}_1 - \hat{\mathcal{T}}_2 - \dots} \hat{\mathbb{H}} e^{\hat{\mathcal{T}}_1 + \hat{\mathcal{T}}_2 + \dots}$$

Cluster-Operators $\hat{\mathcal{T}}_1, \hat{\mathcal{T}}_2, \dots$



Conventional coupled-cluster: Solve for amplitudes

$$t_i^a, t_{ij}^{ab} \in \mathbb{R}$$

Real-space form of coupled-cluster: Solve for cluster-functions¹

$$|\tau_i\rangle \in L^2(\mathbb{R}^3), \quad |\tau_{ij}\rangle \in L^2(\mathbb{R}^6)$$

Projectors to ensure orthogonality to reference states ϕ_i

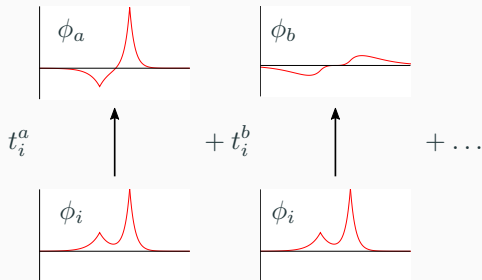
$$Q = 1 - \sum_i |\phi_i\rangle\langle\phi_i|$$

¹R. Bukowski, B. Jeziorski and K. Szalewicz, *J. Chem. Phys.*, 110(9), **1999**.

Fixed-basis formalism

Reference calculation: occupied orbitals
and also virtual orbitals

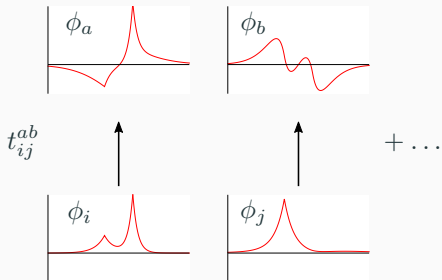
$\hat{\mathcal{T}}_n$ replaces
 n occupied orbitals
by
 n virtual orbitals



Fixed-basis formalism

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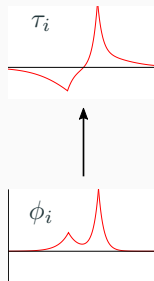
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Real-space formalism

Reference calculation: occupied orbitals

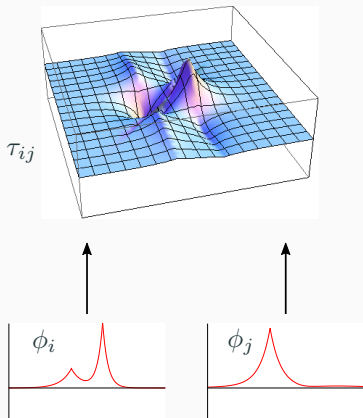
$\hat{\mathcal{T}}_n$ replaces
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Real-space formalism

Reference calculation: occupied orbitals

$\hat{\mathcal{T}}_n$ replaces
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cluster-functions and amplitudes

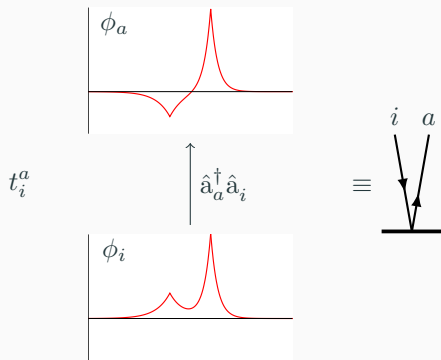
$$|\tau_i\rangle = \sum_a t_i^a |a\rangle$$

$$|\tau_{ij}\rangle = \sum_{ab} t_{ij}^{ab} |ab\rangle$$

⋮

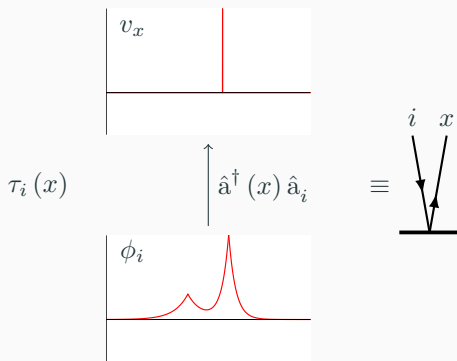
$$|a\rangle = \mathcal{Q}|a\rangle, \quad \langle a|\tau_i\rangle \equiv t_i^a, \quad \langle a|a'\rangle \equiv \delta_{aa'}$$

$$\hat{\mathcal{T}}_1 = \sum_{ia} t_i^a \hat{a}_a^\dagger \hat{a}_i$$



$$|v_x\rangle = \mathcal{Q}|x\rangle, \quad \langle x|\tau\rangle = \tau(x), \quad \langle x'|x\rangle = \delta(x-x')$$

$$\hat{T}_1 = \sum_i \int dx \tau_i(x) \hat{a}^\dagger(x) \hat{a}_i$$



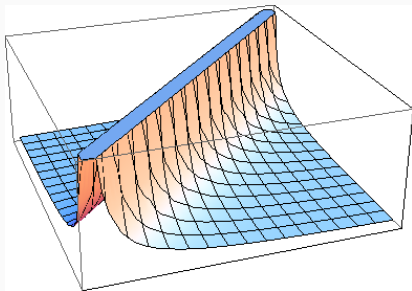
Real-space interpretation of diagrams:

- Only differs for virtual lines
- Summation over virtual indices \rightarrow integration over function variable
- Virtual lines carry \mathcal{Q} projectors

$$\begin{array}{c} i \quad 1 \quad 2 \quad j \\ \downarrow \quad \uparrow \quad \downarrow \quad \downarrow \\ \tau_i \quad \tau_j \end{array} = \int dx'_1 \int dx'_2 \mathcal{Q}_{12} \langle x_1 x_2 | \frac{1}{r_{12}} | x'_2 x'_3 \rangle \tau_i(x'_1) \tau_j(x'_2)$$
$$= \mathcal{Q}_{12} \frac{1}{r_{12}} \tau_i(x_1) \tau_j(x_2)$$

$$r_{12} = \|x_1 - x_2\|$$

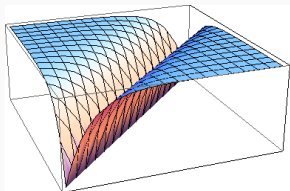
- Electron-repulsion: 3D subspace of 6D space



Need regularization to avoid the direct representation of $\frac{1}{r_{12}}$

$$\left[\hat{T}_{12}, f_{12} \right] = \mathcal{U}_{12} - \frac{1}{r_{12}} \quad \text{cancels } \frac{1}{r_{12}} \text{ in the main equations}$$

$$f_{12} = \frac{1}{2} (1 - e^{-r_{12}})$$



Similar to explicitly correlated methods in LCAO

W. Kutzelnigg, *Theoret. Chim. Acta*, 68(6), 1985.

W. Klopper *et al.*, *Int. Rev. Phys. Chem.*, 25:3, 2006.

What about the nuclear potential?

Regularization is not crucial but it still helps

Key ideas are similar

F. A. Bischoff, *J. Chem. Phys.*, 141, **2014**.

F. F. Seelig, *Z. Naturforsch. A*, 21, **1966**.



Formal scaling N_{occ}^3

Relaxed orbitals and modified projector

$$|t_i\rangle = |i\rangle + |\tau_i\rangle$$

$$Q^t = 1 - \sum_k |t_k\rangle\langle k|$$

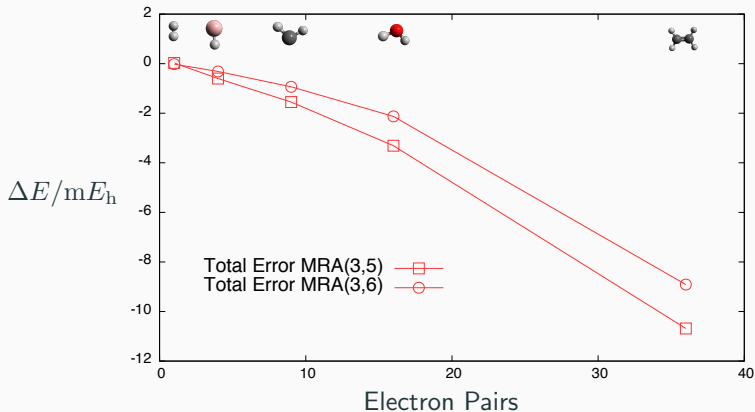
Equivalent to $\hat{\mathcal{T}}_1$ transformed Hamiltonian

Current bottleneck: 6D-MRA for two-body cluster functions

Ground State Correlation Energies



Difference in total correlation energies:
 $\Delta E = \text{MP2-F12/aug-cc-pV6Z} - \text{MP2/MRA}(x,k)$



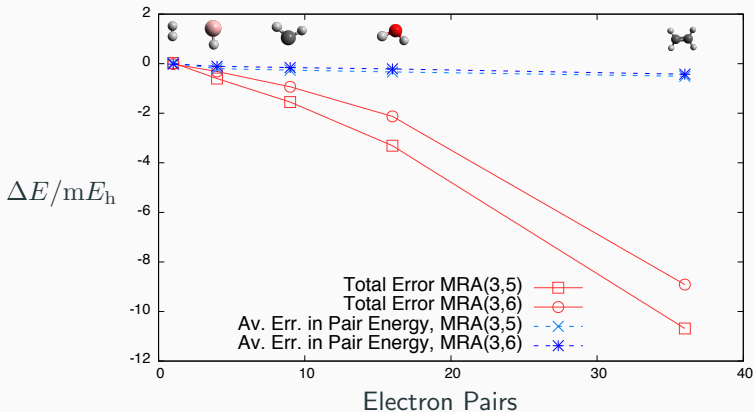
Notation: $\text{MRA}(x,k) \rightarrow \text{MRA threshold of } 10^{-x} \text{ with } k \text{ Legendre scaling functions}$

MP2-F12/aug-cc-pV6Z values obtained with TURBOMOLE

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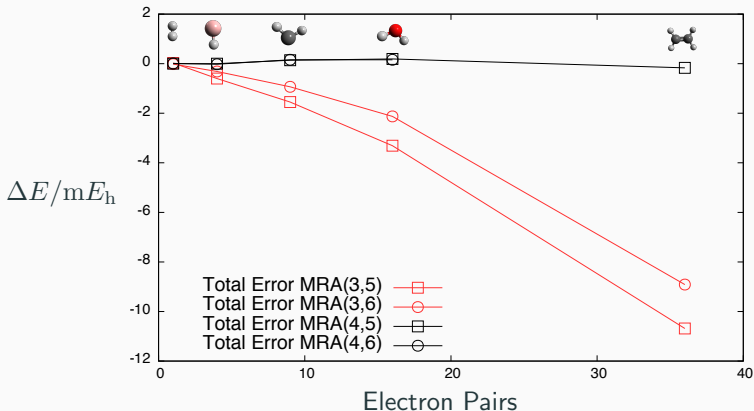
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MP2-F12/aug-cc-pV6Z values obtained with TURBOMOLE

CC2 excitation energies are eigenvalues of the CC2 Jacobian^{1,2}

$$\mathbf{A}\mathbf{x} = \omega\mathbf{x}$$

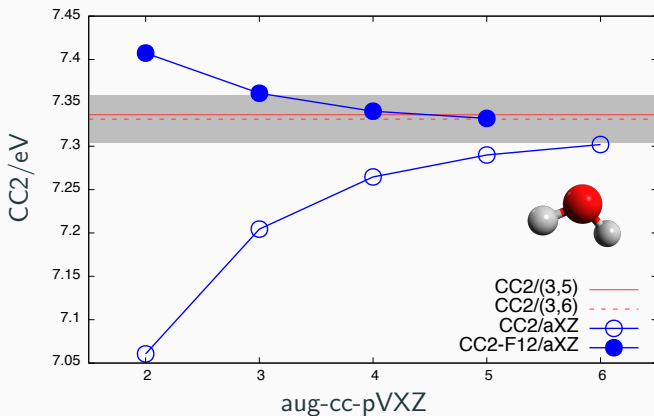
$$\mathbf{A} = \begin{pmatrix} \frac{\partial\Omega_i^a}{\partial t_1} & \frac{\partial\Omega_i^a}{\partial t_2} \\ \frac{\partial\Omega_{ij}^{ab}}{\partial t_1} & \frac{\partial\Omega_{ij}^{ab}}{\partial t_2} \end{pmatrix}$$

Real-space formulation: Take functional derivative of $\Omega[\tau]$

¹O. Christiansen, H. Koch and P. Jørgensen. *Chem Phys. Lett.*, 243(5), **1995**

²H. Koch, P. Jørgensen, *J. Chem. Phys.*, 93, **1990**

Linear Response CC2 - H₂O Example



CC2 with TURBOMOLE, CC2-F12 with KOALA



- Numerically accurate excitation energies for small molecules
- No error accumulation observed
- No balancing issues between ground and excited state
- Accurate representation of one- and two-body terms

PNO-MRA

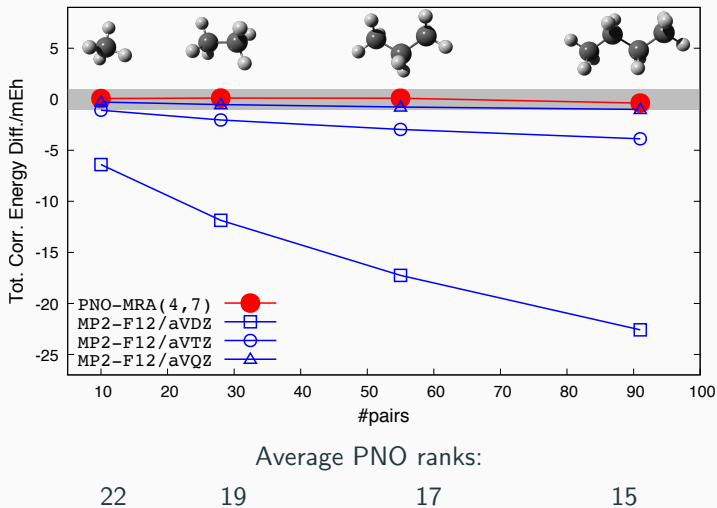


Represent the pair functions by pair natural orbitals (PNOs)

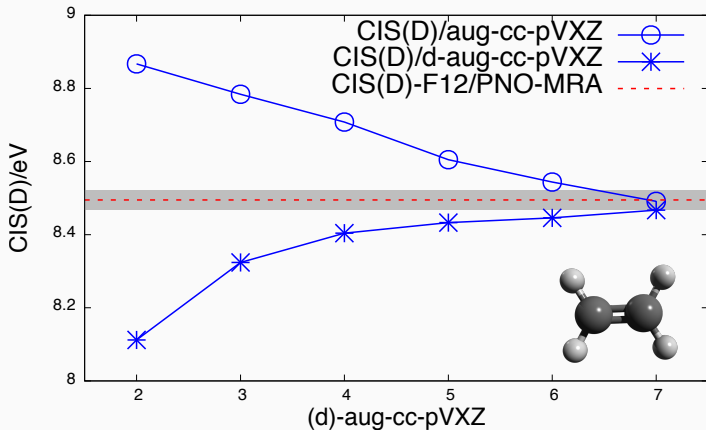
$$|\tau_{ij}\rangle = \sum_{ab} t_{a_{ij}b_{ij}} (|a_{ij}\rangle \otimes |b_{ij}\rangle)$$

Represent and optimize the PNOs with MRA

Difference in total correlation energies: Reference is MP2-F12/aug-cc-pV6Z



Adaptive growing of PNO ranks



Average PNO rank: 43 (canon) or 25 (localized)

Conclusion



MRA-CC2¹

Numerical accuracy is black-box

Exact explicit correlation is crucial

Decreased formal scaling but large prefactor (6D-MRA)

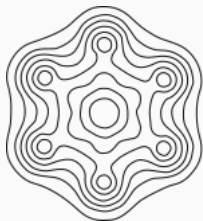
PNO-MRA

3D-MRA for all functions

Low ranks with explicit correlation

Extension to higher excitation levels more straightforward

¹ JSK, F. A. Bischoff, *JCCT*, 13, 2017



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